

# Entanglement, free energy and $C$ -theorem in defect CFT

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based on 1810.06995 with N. Kobayashi, Y. Sato and K. Watanabe  
and a work in progress with K. Goto, L. Nagano and T. Okuda

# Outline

- 1 Defect conformal field theories
- 2 Entanglement entropy and sphere free energy in DCFT
- 3 Towards a  $\mathcal{C}$ -theorem in DCFT

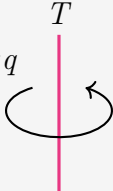
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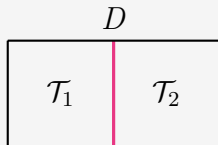
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# Defects in quantum field theory

Defects = Non-local objects in QFTs

- Defined by boundary conditions around them
- Many examples:
  - 1-dim : Line operators (Wilson-'t Hooft loops)
  - 2-dim : Surface operators
- Codim-1 : Domain walls, interfaces and boundaries
- Codim-2 : Entangling surface for entanglement entropy

$$\int_{\mathbb{S}^2} F = 2\pi q$$




## Defect as a probe of QFT phases

- Characterize phases of QFTs: [’t Hooft 78]

$W$ : Wilson loop,  $T$ : ’t Hooft loop

Confinement :  $\langle W \rangle \sim e^{-\text{Area}}$ ,  $\langle T \rangle \sim e^{-\text{Length}}$

Higgs :  $\langle W \rangle \sim e^{-\text{Length}}$ ,  $\langle T \rangle \sim e^{-\text{Area}}$

- Higher-dimensional generalization:
  - Wilson surface operators,

$$W_{\Sigma} = \exp \left( i \int_{\Sigma} A \right)$$

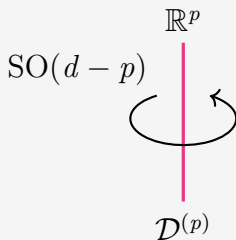
for a  $p$ -form gauge field  $A$

# Conformal defects

- In Euclidean CFT $_d$ , the conformal group is  $SO(d + 1, 1)$
- $p$ -dimensional conformal defects  $\mathcal{D}^{(p)}$  are either **flat** or **spherical**, preserving

$SO(p + 1, 1)$  : conformal symmetry on defects

$SO(d - p)$  : rotation in the transverse direction



- Defects allow for **defect local operators**  $\hat{\mathcal{O}}_n(\hat{x})$   
 $\hat{x}^a$ : parallel coordinates on  $\mathcal{D}^{(p)}$  ( $a = 1, \dots, p$ )

# One-point function

The residual conformal symmetry constrains one-point functions

- Scalar primary:

$$\langle \mathcal{O}(x) \rangle^{(\text{DCFT})} = \frac{a_{\mathcal{O}}}{|x_{\perp}|^{\Delta}}, \quad x_{\perp}^i : \text{transverse coordinates}$$

$$(i = p + 1, \dots, d)$$

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- Stress tensor:

$$\langle T^{ab}(x) \rangle^{(\text{DCFT})} = \frac{d - p - 1}{d} \frac{a_T}{|x_{\perp}|^d} \delta^{ab}$$

$$\langle T^{ij}(x) \rangle^{(\text{DCFT})} = -\frac{a_T}{|x_{\perp}|^d} \left( \frac{p + 1}{d} \delta^{ij} - \frac{x_{\perp}^i x_{\perp}^j}{|x_{\perp}|^2} \right)$$

$$\langle T^{ai}(x) \rangle^{(\text{DCFT})} = 0$$



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$$\langle T^{ai}(x) \rangle^{(\text{DCFT})} = 0$$

N.B.  $\langle T^{\mu\nu}(x) \rangle^{(\text{DCFT})} = 0$  in BCFT ( $p = d - 1$ ) [McAvity-Osborn 95]

# Applications of conformal defects

## ■ Constrain bulk CFT data in defect CFT by crossing symmetry

[Liendo-Rastelli-van Rees 12, Gaiotto-Mazac-Paulos 13, Gliozzi-Liendo-Meineri-Rago 15, Billó-Gonçalves-Lauria-Meineri 16, Lemos-Liendo-Meineri-Sarkar 17, ...]

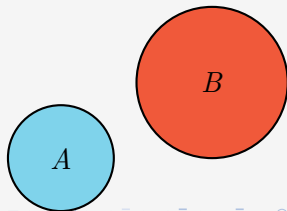
$$\sum_k \begin{array}{c} \mathcal{O}_1 \quad \mathcal{O}_2 \\ \diagdown \quad \diagup \\ \text{Y} \\ \diagup \quad \diagdown \\ \mathcal{O}_k \\ | \\ \hline \mathcal{D} \end{array} = \sum_l \begin{array}{c} \mathcal{O}_1 \quad \mathcal{O}_2 \\ | \quad | \\ \hline \hat{\mathcal{O}}_l \\ | \quad | \\ \hline \mathcal{D} \end{array}$$

## ■ Understand quantum entanglement in QFT:

e.g. **Mutual information** as a correlator of two defects

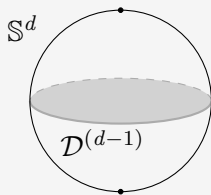
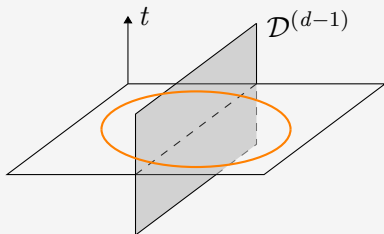
[Cardy 13, Bianchi-Meineri-Myers-Smolkin 15, Chen-Chen-Hao-Long 17]

$$\begin{aligned} I(A, B) &\equiv S_A + S_B - S_{A \cup B} \\ &= \log \frac{\langle \mathcal{D}(\partial A) \mathcal{D}(\partial B) \rangle}{\langle \mathcal{D}(\partial A) \rangle \langle \mathcal{D}(\partial B) \rangle} \end{aligned}$$



# Goal of this talk

- In DCFT we will study
  - Entanglement entropy across a sphere
  - Sphere free energy
- How do they depend on defect data?
- What is the measure of degrees of freedom associated to defects?



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## Definition of entanglement entropy

Divide a system to  $A$  and  $B = \bar{A}$ :  $\mathcal{H}_{\text{tot}} = \mathcal{H}_A \otimes \mathcal{H}_B$

### Entanglement entropy

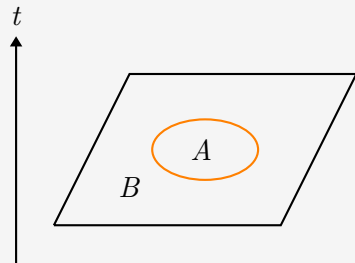
$$S_{\text{EE}} = -\text{tr}_A \rho_A \log \rho_A$$

- The reduced density matrix

$$\rho_A \equiv \text{tr}_B \rho_{\text{tot}}$$

- For a pure ground state  $|\Psi\rangle$

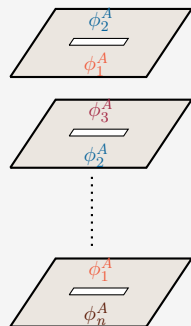
$$\rho_{\text{tot}} = |\Psi\rangle \langle \Psi|$$



# Replica trick and Rényi entropy

Entanglement entropy

$$S_{\text{EE}} = \lim_{n \rightarrow 1} S_n$$



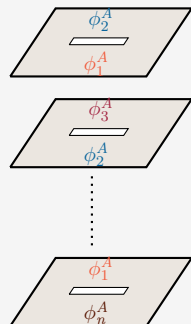
# Replica trick and Rényi entropy

## Entanglement entropy

$$S_{\text{EE}} = \lim_{n \rightarrow 1} S_n$$

## $n^{\text{th}}$ Rényi entropy

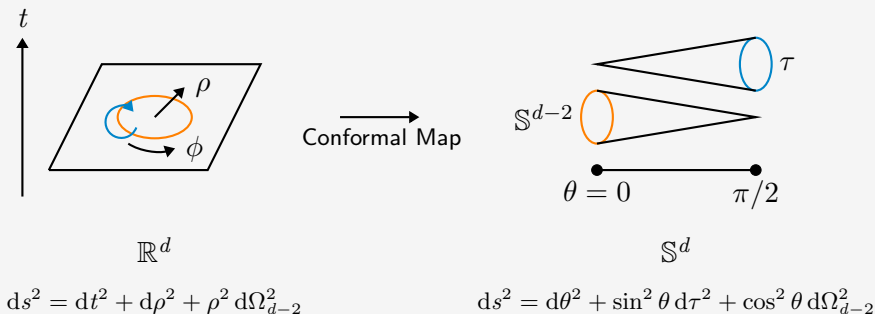
$$S_n = \frac{1}{1-n} \log \frac{Z_n}{(Z_1)^n}$$



$Z_n$ : partition function on the  $n$ -fold cover branched over  $A$

## Conformal map for a spherical region

- Exact calculations limited to a few simple cases (free fields, planar  $\partial A$  in CFT, etc)
- For a spherical region in CFT, however, there exists a **conformal map** to the  $n$ -fold cover of  $\mathbb{S}^d$  (**CHM map**) [Casini-Huerta-Myers 11]





## Entanglement entropy across a sphere

- CFT partition function is invariant under the CHM map

$$Z_n^{(\text{CFT})} = Z^{(\text{CFT})}[\mathbb{S}_n^d], \quad \mathbb{S}_n^d: n\text{-fold cover of } \mathbb{S}^d \quad (\tau \sim \tau + 2\pi n)$$

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- Hence the Rényi entropy is given by

The Rényi entropy across a sphere in CFT

$$S_n = \frac{1}{1-n} \log \frac{Z^{(\text{CFT})}[\mathbb{S}_n^d]}{(Z^{(\text{CFT})}[\mathbb{S}^d])^n}$$

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- Taking  $n \rightarrow 1$  limit

Entanglement entropy across a sphere in CFT [Casini-Huerta-Myers 11]

$$S_{\text{EE}} = \log Z^{(\text{CFT})}[\mathbb{S}^d]$$

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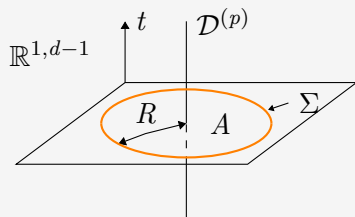
Entanglement entropy across a sphere in CFT [Casini-Huerta-Myers 11]

$$S_{\text{EE}} = \log Z^{(\text{CFT})}[\mathbb{S}^d]$$

- For free fields,  $Z^{(\text{CFT})}[\mathbb{S}_n^d]$ : one-loop determinant [Klebanov-Pufu-Sachdev-Safdi 11]
- For SUSY gauge theories,  $Z^{(\text{CFT})}[\mathbb{S}^d]$  ( $n = 1$ ) can be calculated by **SUSY localization** [Pestun 07, Kapustin-Willet-Yaakov 09, ...]

# Entanglement entropy across a sphere with defects

- $A$ : a ball centered at the origin
- $\mathcal{D}^{(p)}$ : a  $p$ -dim flat defect
- After a conformal transformation  
(**CHM map** [Casini-Huerta-Myers 11, Jensen-O'Bannon 13])



The  $n$ -th Rényi entropy

$$S_n^{(\text{DCFT})} = \frac{1}{1-n} \log \frac{Z^{(\text{DCFT})}[\mathbb{S}_n^d]}{(Z^{(\text{DCFT})}[\mathbb{S}^d])^n}$$

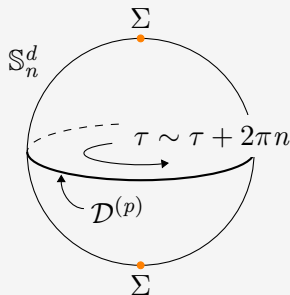
$\mathbb{S}_n^d$ :  $n$ -fold cover of  $\mathbb{S}^d$

# Defect entropy

- The excess of EE is measured by

## Defect entropy

$$S_{\text{defect}} \equiv \lim_{n \rightarrow 1} \left( S_n^{(\text{DCFT})} - S_n^{(\text{CFT})} \right)$$



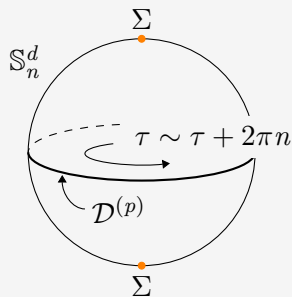
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$$\begin{aligned}
 S_{\text{defect}} &\equiv \lim_{n \rightarrow 1} \left( S_n^{(\text{DCFT})} - S_n^{(\text{CFT})} \right) \\
 &= \lim_{n \rightarrow 1} \frac{1}{1-n} \log \frac{\langle \mathcal{D}^{(p)} \rangle_n}{\langle \mathcal{D}^{(p)} \rangle_1}
 \end{aligned}$$

$$\begin{aligned}
 \langle \mathcal{D}^{(p)} \rangle_n &\equiv \frac{Z^{(\text{DCFT})}[\mathbb{S}_n^d]}{Z^{(\text{CFT})}[\mathbb{S}_n^d]} \\
 \langle \mathcal{D}^{(p)} \rangle &\equiv \langle \mathcal{D}^{(p)} \rangle_1 \quad (\text{vev of } \mathcal{D}^{(p)})
 \end{aligned}$$



$n \rightarrow 1$  limit

- Expansion around  $n = 1$  ( $\delta g_{\mu\nu} = O(n - 1)$ )

$$\log Z^{(\text{DCFT})}[\mathbb{S}_n^d] = \log Z^{(\text{DCFT})}[\mathbb{S}^d] - \frac{1}{2} \int_{\mathbb{S}^d} \delta g_{\mu\nu} \langle T^{\mu\nu} \rangle_{\mathbb{S}^d}^{(\text{DCFT})} + O((n - 1)^2)$$



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- In CFT [Casini-Huerta-Myers 11]

$$\langle T^{\mu\nu} \rangle_{\mathbb{S}^d}^{(\text{CFT})} = 0$$

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- In DCFT

$$\langle T^{\mu\nu} \rangle_{\mathbb{S}^d}^{(\text{DCFT})} \neq 0 \quad (\propto a_T)$$

# Universal relation

Defect entropy and sphere free energy [Kobayashi-TN-Sato-Watanabe 18]

$$S_{\text{defect}} = \log \langle \mathcal{D}^{(p)} \rangle - \frac{2(d-p-1)\pi^{d/2+1}}{\sin(\pi p/2) d \Gamma(p/2+1) \Gamma((d-p)/2)} a_T$$

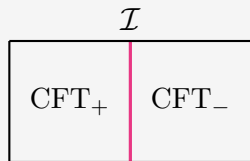
- Dimensional regularization is assumed
- Equality holds up to UV divergences
- Reproduces a known result when  $p = 1$  [Lewkowycz-Maldacena 13]
- The second term in rhs vanishes when  $p = d - 1$  (codimension-one)

# Interface entropy

- Interface CFT,  $\mathcal{D}^{(d-1)} = \mathcal{I}$

Interface entropy

$$S_{\mathcal{I}} = S^{(\text{ICFT})} - \frac{S^{(\text{CFT}_+)} + S^{(\text{CFT}_-)}}{2}$$



- Universal relation:

$$S_{\mathcal{I}} = \log \langle \mathcal{I} \rangle, \quad \langle \mathcal{I} \rangle \equiv \frac{Z^{(\text{ICFT})}[\mathbb{S}^d]}{(Z^{(\text{CFT}_+)}[\mathbb{S}^d] Z^{(\text{CFT}_-)}[\mathbb{S}^d])^{1/2}}$$

## Half-BPS Janus interfaces

- In  $2d$   $\mathcal{N} = (2, 2)$  and  $4d$   $\mathcal{N} = 2$  SCFTs

$$Z^{(\text{SCFT})}[\mathbb{S}^d](\tau, \bar{\tau}) = \left(\frac{r}{\epsilon}\right)^{-4A} \exp[K(\tau, \bar{\tau})/12]$$

- $r$ : radius,  $\epsilon$ : UV cutoff,  $A$ : type- $A$  central charge
- $K(\tau, \bar{\tau})$ : Kähler potential on a conformal manifold

[Jockers-Kumar-Lapan-Morrison-Romo 12, Gomis-Lee 12, Gerchkovitz-Gomis-Komargodski 14]

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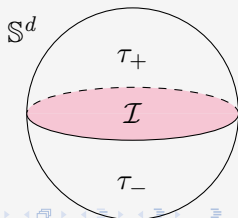
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- Half-BPS Janus interfaces [Drukker-Gaiotto-Gomis 10, Goto-Okuda 18]

$$Z^{(\text{ICFT})}[\mathbb{S}^d] = Z^{(\text{SCFT})}[\mathbb{S}^d](\tau_+, \bar{\tau}_-)$$

$\exists$  Kähler ambiguity [Gomis-Ishtiaque 14]

$$K(\tau_+, \bar{\tau}_-) \rightarrow K(\tau_+, \bar{\tau}_-) + \mathcal{F}(\tau_+) + \bar{\mathcal{F}}(\bar{\tau}_-)$$



# Interface entropy as Calabi's diastasis

Using supersymmetric Rényi entropy [TN-Yaakov 13] to preserve SUSY

Interface entropy as Calabi's diastasis [Goto-Nagano-TN-Okuda, WIP]

$$S_{\mathcal{I}} = \log |\langle \mathcal{I} \rangle|$$



## Interface entropy as Calabi's diastasis

Using supersymmetric Rényi entropy [TN-Yaakov 13] to preserve SUSY

Interface entropy as Calabi's diastasis [Goto-Nagano-TN-Okuda, WIP]

$$S_{\mathcal{I}} = -\frac{1}{24} [K(\tau_+, \bar{\tau}_+) + K(\tau_-, \bar{\tau}_-) - K(\tau_+, \bar{\tau}_-) - K(\tau_-, \bar{\tau}_+)]$$

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- Bulk part ( $\propto A$ ) cancels out
- No Kähler ambiguity
- Agree with the known result in  $2d$  [Bachas-BrunnerDouglas-Rastelli 13, Bachas-Plencner 16]
- Reproduces a conjectured form in  $4d$  [Goto-Okuda 18]

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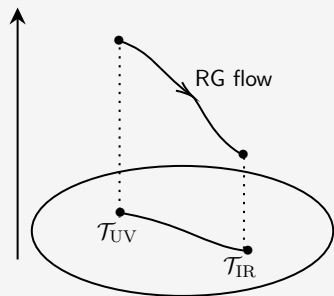
# $\mathcal{C}$ -theorem

## $\mathcal{C}$ -theorem (weak)

$\exists$  a function  $\mathcal{C}(\mathcal{T})$  on a theory space  
s.t.

$$\begin{aligned} \mathcal{T}_{\text{UV}} &\xrightarrow{\text{RG flow}} \mathcal{T}_{\text{IR}} \\ \Rightarrow \quad \mathcal{C}(\mathcal{T}_{\text{UV}}) &\geq \mathcal{C}(\mathcal{T}_{\text{IR}}) \end{aligned}$$

$\mathcal{C}(\mathcal{T})$  (height function on  $\mathcal{S}$ )



$\mathcal{S}$  = space of QFTs

- $\mathcal{C}(\mathcal{T})$  called a  $\mathcal{C}$ -function ( $\approx$  resource measure)
- Regarded as a **measure of degrees of freedom** in QFT
- Constrains the dynamics under RG if holds

# Examples and conjectures

- $2d$ : Zamolodchikov's  $c$ -theorem [Zamolodchikov 86]
- even  $d$ :  $A$ -theorem ( $\langle T_{\mu}^{\mu} \rangle_{\mathbb{S}^d} \propto A$ ) [Cardy 88, Komargodski-Schwimmer 11]

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Generalized  $F$ -theorem conjecture [Giombi-Klebanov 14]

$$\tilde{F} \equiv \sin\left(\frac{\pi d}{2}\right) \log Z[\mathbb{S}^d], \quad \tilde{F}_{\text{UV}} \geq \tilde{F}_{\text{IR}}$$



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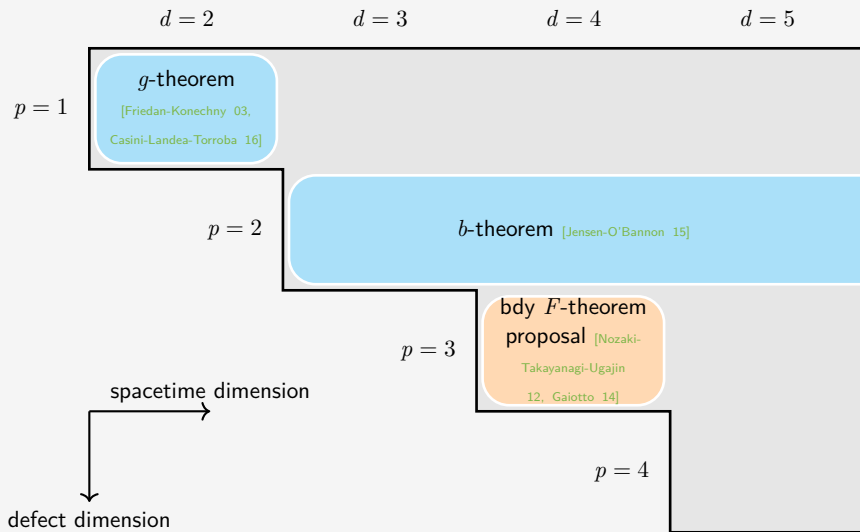
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- Reduces to the  $F$ - and  $A$ -theorems

$$\tilde{F} = \begin{cases} F & d : \text{odd} \\ \frac{\pi}{2} A \text{ (conformal anomaly)} & d : \text{even} \end{cases}$$

Status of  $\mathcal{C}$ -theorems (+ conjectures) in DCFT

# $\mathcal{C}$ -function in DCFT?

- Candidate of  $\mathcal{C}$ -functions
  - **entanglement entropy**: holographic model ( $p = d - 1$ )  
[Estes-Jensen-O'Bannon-Tsatis-Wrase 14]
  - **sphere free energy**: bdy  $F$ -thm,  $b$ -thm ( $p = 2$ ),  
Wilson loop RG flow ( $d = 4, p = 1$ ) [Beccaria-Giombi-Tseytlin 17]
- These two agree when  $p = d - 1$  due to the universal relation

Are both  $\mathcal{C}$ -functions for any  $d$  and  $p$ ?

## $\mathcal{C}$ -theorem in DCFT: conjecture

- Defect RG flow triggered by a relevant defect operator:

$$I = I_{\text{DCFT}} + \hat{\lambda} \int d^p \hat{x} \sqrt{\hat{g}} \hat{\mathcal{O}}(\hat{x})$$

Conjecture [Kobayashi-TN-Sato-Watanabe 18]

The universal part of the sphere free energy

$$\tilde{D} \equiv \sin\left(\frac{\pi p}{2}\right) \log |\langle \mathcal{D}^{(p)} \rangle|$$

does not increase along any defect RG flow

$$\tilde{D}_{\text{UV}} \geq \tilde{D}_{\text{IR}}$$

- Same form as the generalized  $F$ -thm:  $\tilde{F} = \sin\left(\frac{\pi d}{2}\right) \log Z[\mathbb{S}^d]$

# Checks

- Sphere free energy always decreases under defect RG flows in
  - Conformal perturbation theory
  - Wilson loops ( $p = 1$ ) in  $3d$  and  $4d$
  - Holographic models (a proof assuming null energy condition)
  
- Entanglement entropy does not decrease in
  - Wilson loop RG flows [Kobayashi-TN-Sato-Watanabe 18], surface operators ( $p = 2$ ) [Jensen-O'Bannon-Robinson-Rodgers 18]
  - Holographic Wilson loops [Kumar-Silvani 16, 17] and surface operators [Rodgers 18]

$$S_{\text{defect}} = \log \langle \mathcal{D}^{(p)} \rangle - \frac{2(d-p-1)\pi^{d/2+1}}{\sin(\pi p/2) d \Gamma(p/2+1) \Gamma((d-p)/2)} a_T$$

# Summary and future work

## ■ Summary:

- Find the universal relation between EE and sphere free energy
- Derive the interface entropy as Calabi's diastasis
- Propose a  $\mathcal{C}$ -theorem in DCFT

## ■ Future work:

- Proof in SUSY theories? (cf.  $F$ - and  $a$ -maximizations [Jafferis 10, Closset-Dumitrescu-Festuccia-Komargodskia-Seiberg 12, Intriligator-Wecht 03])
- Proof using entropic inequalities as in  $g$ -thm [Casini-Landea-Torroba 16]?
- Constrains on the dynamics of defect RG flows?
- Dependence of defect entropy on bulk marginal deformation? (cf. [Herzog-Shamir 18, Bianchi 18] for sphere free energy)